

CHEBYSHEV NETS IN A SUBSPACE OF A GENERALIZED WEYL SPACE

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ABSTRACT

Let GW_m be m dimensional subspace of an n -dimensional subspace of an n -dimensional generalized Weyl space GW_n and let $\delta = (v_1, v_2, \dots, v_m)$ be a net in GW_m . In this paper the necessary and the sufficient condition that the net δ is to be a chebyshev net and a geodesic net relative to GW_m and GW_n is obtained.

Keywords: Generalized Weyl Space, Chebyshev Net, Geodesic Net.

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GENELLEŞTİRİLMİŞ WEYL UZAYININ ALT UZAYLARINDA CHEBYSHEV ŞEBEKELERİ

ÖZET

n boyutlu GW_n genelleştirilmiş Weyl uzayının m boyutlu GW_m alt uzayına ait bir şebeke $\delta = (v_1, v_2, \dots, v_m)$ olsun. Bu çalışmada δ şebekesinin GW_m ve GW_n e göre Chebyshev şebekesi ve jeodezik şebeke olması için gerek ve yeter koşullar elde edilmiştir.

Anahtar Sözcükler: Genelleştirilmiş weyl uzayı, Chebyshev şebekesi Geodesic şebeke.

1. INTRODUCTION

An m dimensional GW_m is said to be generalized Weyl space if it has an asymmetric conformal metric tensor g_{ij} and an asymmetric connection ∇_k satisfying the compatibility condition given by equation

$$\nabla_k g_{ij} = 2T_k g_{ij} \quad (1.1)$$

where T_k denotes a covariant vector field and ∇_k denotes the usual covariant derivative

Under a renormalization of the fundamental tensor of the form $\check{g}_{ij} = \lambda^2 g_{ij}$, the complementary vector field T_k is transformed by the law $\check{T}_k = T_k + \partial_k \ln \lambda$, where λ is a scalar function on GW_m .

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Let L_{jk}^i be denote the coefficients of the asymmetric connection ∇_k . So, a generalized Weyl space is shortly written as $GW_m(L_{jk}^i, g_{ij}, T_k)$.

The main properties of $GW_m(L_{jk}^i, g_{ij}, T_k)$ can be expressed as follows

$$g_{ij} = g_{(ij)} + g_{[ij]} \quad (1.2)$$

$$\nabla_k g_{(ij)} = 2 g_{(ij)} T_k \quad (1.3)$$

$$\nabla_k g_{[ij]} = 2 g_{[ij]} T_k \quad (1.4)$$

$$g_{(ik)} g^{(kl)} = \delta_i^l \quad (1.5)$$

$$\nabla_k g^{(ij)} = -2 T_k g^{(ij)} \quad (1.6)$$

where $g_{(ij)}$ and $g_{[ij]}$ denote symmetric and antisymmetric parts of g_{ij} respectively.

The symmetric part of the connection coefficients L_{jk}^i are given as ([1], [2], [3], [4])

$$L_{(jk)}^i = W_{jk}^i = \begin{bmatrix} i \\ jk \end{bmatrix} - (\delta_j^i T_k + \delta_k^i T_j - g_{(jk)} g^{(mi)} T_m) \quad (1.7)$$

where $\begin{bmatrix} i \\ jk \end{bmatrix}$ are second kind Christoffel symbols defined by

$$\begin{bmatrix} i \\ jk \end{bmatrix} = \frac{1}{2} g^{(ir)} \left[\frac{\partial g_{(jr)}}{\partial x^k} + \frac{\partial g_{(kr)}}{\partial x^j} - \frac{\partial g_{(jk)}}{\partial x^r} \right] \quad (1.8)$$

A quantity A is called a satellite of weight $\{p\}$ of tensor g_{ij} , if it admits a transformation of the form

$$\tilde{A} = \lambda^p A \quad (1.9)$$

The prolonged covariant derivative of a satellite A of the tensor g_{ij} of weight $\{p\}$ is defined by

$$\dot{\nabla}_k A = \nabla_k A - p T_k A \quad (1.10)$$

Let $GW_m(L_{jk}^i, g_{ij}, T_k)$ be a subspace, with coordinates u^i , of the Weyl space $GW_n(L_{\beta\gamma}^\alpha, g_{\alpha\beta}, T_\gamma)$ with coordinates x^α . Suppose that the metrics of $GW_m(L_{jk}^i, g_{ij}, T_k)$ and $GW_n(L_{\beta\gamma}^\alpha, g_{\alpha\beta}, T_\gamma)$ are elliptic and that they are given, respectively, by $g_{ij} du^i du^j$ and $g_{\alpha\beta} du^\alpha du^\beta$ which are connected by the relations

$$g_{ij} = g_{\alpha\beta} x_i^\alpha x_j^\beta \quad (1.11)$$

where x_i^α denotes the covariant derivative of x^α with respect to u^i .

Let v_r^i ($r = 1, 2, \dots, m$) be the contravariant components of the m independent vector fields v_r in GW_m which are normalized by the condition $g_{ij} v_r^i v_r^j = 1$. The vector fields v_1, v_2, \dots, v_m determine a net (v_1, v_2, \dots, v_m) in GW_m .

Corresponding to the vector fields v_r ; the covector fields v^r are defined by the conditions

$$v_r^i v_j^r = \delta_j^i, \quad v_r^i v_i^p = \delta_r^p \tag{1.12}$$

Denote by n_σ^α ($1, 2, \dots, n - m$), the contravariant components of the $n - m$ linear independent unit vector fields n_σ in GW_n normal to GW_m .

The moving frame $\{x_\alpha^i, n_\sigma^\alpha\}$ on GW_m , reciprocal to the moving frame $\{x_i^\alpha, n_\sigma^\alpha\}$ is defined by the relations

$$n_\alpha^\mu n_\sigma^\alpha = \delta_\sigma^\mu, \quad n_\alpha^\sigma x_i^\alpha = 0, \quad x_\alpha^i n_\sigma^\alpha = 0, \quad x_i^\alpha x_\alpha^j = \delta_i^j \tag{1.13}$$

($\mu = 1, 2, \dots, n - m$)

On the other hand, if the components of v_r and v^r relative to GW_n , are respectively denoted by v_r^α and v_α^r , we have

$$v_r^\alpha = v_r^i x_i^\alpha, \quad v_\alpha^r = v_i^r x_\alpha^i \tag{1.14}$$

The derivation formula for v_r and v^r are given by [3], [5]

$$\dot{\nabla}_k v_r^i = T_k^p v_r^i, \quad \dot{\nabla}_k v_i^r = -T_k^p v_i^r \tag{1.15}$$

Taking the prolonged covariant derivative of (1.13)₄ with respect to u^k and remembering that [5]

$$\dot{\nabla}_k x_i^\alpha = \sum_{\sigma=1}^{n-m} w_{ik}^\sigma n_\sigma^\alpha + A_{ik}^h x_h^\alpha \tag{1.16}$$

where w_{ik}^σ are the components of second fundamental form related to the normal n^σ of GW_m defined by

$$w_{ik}^\sigma = n_\alpha^\sigma \dot{\nabla}_k x_i^\alpha \tag{1.17}$$

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and A_{ik}^h are defined by

$$A_{ik}^h = x_\alpha^h \dot{\nabla}_k x_i^\alpha \quad (1.18)$$

From (1.16)₄, we have

$$\nabla_k x_\alpha^j = \sum_{\sigma=1}^{n-m} \Omega_{\sigma k}^j n_\alpha^\sigma - A_{hk}^j x_\alpha^h, \quad \Omega_{\sigma k}^i = n^\alpha \nabla_k x_\alpha^j \quad (1.19)$$

For the definition of the Chebyshev nets of the first and second kind, and that the b-nets and c-nets, the reader is referred to [3].

2. CHEBYSHEV NETS IN A GENERALIZED WEYL SUBSPACE

Consider the net (v_1, v_2, \dots, v_m) in GW_m determined by the vector fields v_1, v_2, \dots, v_m . Let \bar{a}_{rp}^α , a_{rp}^i ; \bar{b}_α^r , b_i^r and \bar{c}_r^α , c_r^i ; be the covariant components of the chebyshev vector fields of the first and the second kind and the geodesic vector fields of the given net relative to GW_n and GW_m , respectively.

Theorem 2.1. For the Chebyshev vector fields of the first and second kind and the geodesic vector fields of the net (v_1, v_2, \dots, v_m) in GW_m we have

$$(i) \quad \bar{a}_{rp}^\alpha = v_p^\gamma \dot{\nabla}_\gamma A = \sum_{\sigma=1}^{n-m} w_{ik} v_p^k v_r^i n_\sigma^\alpha + a_{rp}^h x_\alpha^h + A_{ik}^h v_r^i v_p^k x_\alpha^h$$

$$w_{ik}^\sigma = g_{(\alpha\beta)} n_\sigma^\beta \dot{\nabla}_k x_i^\alpha \quad (r \neq p)$$

$$(ii) \quad \bar{b}_\alpha^r = v_p^\gamma \dot{\nabla}_\gamma v_\alpha^r = - \sum_{\sigma=1}^{n-m} \Omega_{\sigma k}^i v_i^r v_p^k n_\alpha^\sigma + b_h^r x_\alpha^h + A_{hk}^i x_\alpha^h v_i^r v_p^k$$

$$\Omega_{\sigma k}^j = n^\alpha \dot{\nabla}_k x_\alpha^j$$

$$(iii) \quad \bar{c}_r^\alpha = v_r^\gamma \dot{\nabla}_\gamma v_r^\alpha = \sum_{\sigma=1}^{n-m} w_{ik} v_r^i v_r^k n_\sigma^\alpha + x_h^\alpha c_r^h + A_{ik}^h v_r^i v_r^k x_h^\alpha$$

where the symbol $\textcircled{\circ}$ indicates that the summation on R will not be made.

Proof. (i) Taking the prolonged covariant derivative of each side of (1.14)₁ with respect to u^k and using (1.15) and (1.16), we get

$$\dot{\nabla}_k v_r^\alpha = v_r^i (\dot{\nabla}_k x_i^\alpha) + x_i^\alpha (\dot{\nabla}_k v_r^i) = v_r^i \left(\sum_{\sigma=1}^{n-m} w_{ik} n_\sigma^\alpha + A_{ik}^h x_h^\alpha \right) + x_i^\alpha \left(T_{kp}^i v_p^i \right) \quad (2.1)$$

$$v_l^k \overset{\bullet}{\nabla}_k v_r^\alpha = \sum_{\sigma=1}^{n-m} w_{ik} v_r^i v_l^k n_\sigma^\alpha + A_{ik}^h v_r^i v_l^k x_h^\alpha + x_i^\alpha (T_k^p v_r^k) v_l^i \quad (2.2)$$

In view of

$$v_l^k \overset{\bullet}{\nabla}_k = v_l^\gamma \overset{\bullet}{\nabla}_\gamma, \quad T_k^p v_l^k = \tau_{rl}^p \quad (2.3)$$

the last relation becomes

$$v_l^k \overset{\bullet}{\nabla}_\gamma v_r^\alpha = \sum_{\sigma=1}^{n-m} w_{ik} v_r^i v_l^k n_\sigma^\alpha + A_{ik}^h v_r^i v_l^k x_h^\alpha + x_i^\alpha (\tau_{rl}^p v_l^i) \quad (2.4)$$

where the functions τ_{rl}^p ($r \neq l$) are the chebyshev curvatures of the first kind of the curves of the net

$$a_{rp} = \tau_{rp}^s v_{rp}^i \quad (2.5)$$

So that (2.4) reduces to

$$\bar{a}_{rl}^\alpha = v_l^\gamma \overset{\bullet}{\nabla}_\gamma A = \sum_{\sigma=1}^{n-m} w_{ik} v_l^k v_r^i n_\sigma^\alpha + a_{rl}^h x_h^\alpha + A_{ik}^h v_r^i v_l^k x_h^\alpha \quad (r \neq l) \quad (2.6)$$

(ii) Differentiating covariantly both sides of (1.14)₂ with respect to u^k and taking (1.19) and (1.15)₂ into consideration we obtain

$$\begin{aligned} \overset{\bullet}{\nabla}_k v_\alpha^r &= \overset{\bullet}{\nabla}_k (v_i x_\alpha^i) = (\overset{\bullet}{\nabla}_k v_i) x_\alpha^i + v_i (\overset{\bullet}{\nabla}_k x_\alpha^i) \\ &= (\sum_{\sigma=1}^{n-m} \Omega_k^i n_\sigma^\alpha - A_{hk}^i x_\alpha^h) v_i - T_k^p v_i x_\alpha^i \end{aligned} \quad (2.7)$$

Multiplying (2.7) by v_α^k and remembering that $\overset{\bullet}{\nabla}_k v_\alpha^r = x_k^\gamma \overset{\bullet}{\nabla}_\gamma v_\alpha^r$, we have

$$v_\alpha^k \overset{\bullet}{\nabla}_\gamma v_\alpha^r = \sum_{\sigma=1}^{n-m} \Omega_k^i n_\sigma^\alpha v_i v_\alpha^k - A_{hk}^i v_i v_\alpha^k x_\alpha^h - T_k^p v_i v_\alpha^k x_\alpha^i \quad (2.8)$$

Since the Chebyshev curvatures ρ_l^r of the second kind of the curves belonging to the net (v_1, v_2, \dots, v_m) are defined by

$$\rho_l^r = T_k^p v_\alpha^k \quad (2.9)$$

The last relation takes the form

$$v_\alpha^k \overset{\bullet}{\nabla}_\gamma v_\alpha^r = \sum_{\sigma=1}^{n-m} \Omega_k^i n_\sigma^\alpha v_i v_\alpha^k - A_{hk}^i v_i v_\alpha^k x_\alpha^h - (\rho_l^r v_i) x_\alpha^i \quad (2.10)$$

or, in terms of the chebyshev vector fields of the second kind, (2.10) becomes

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$$\frac{r}{b} \alpha = - \sum_{\sigma=1}^{n-m} \Omega_k^i \underset{\textcircled{R}}{v_i} v^k \underset{\textcircled{R}}{n_\alpha} - b_h^r x_\alpha^h - A_{hk}^i x_\alpha^h \underset{\textcircled{R}}{v_i} v^k \quad (2.11)$$

where $\overset{r}{b}_i$ is defined by $\overset{r}{b}_i = \underset{s}{\rho} v_i^s$.

(iii) If we multiply (2.1) by $\underset{r}{v}^k$ and sum for k , we find

$$\underset{r}{v}^k \overset{\bullet}{\nabla}_k \underset{r}{v}^\alpha = \sum_{\sigma=1}^{n-m} w_{ik} v^i v^k n^\alpha + (T_k^p v^k) v^i x_i^\alpha + A_{ik}^h v^i v^k x_h^\alpha \quad (2.12)$$

or using (1.14)₁ and the fact that $\overset{\bullet}{\nabla}_k = x_k^\gamma \overset{\bullet}{\nabla}_\gamma$, we transform the last equation into

$$\underset{r}{v}^\gamma \overset{\bullet}{\nabla}_\gamma \underset{r}{v}^\alpha = \sum_{\sigma=1}^{n-m} w_{ik} v^i v^k n^\alpha + \overset{p}{S} v^i x_i^\alpha + A_{ik}^h v^i v^k x_h^\alpha \quad (2.13)$$

where $\overset{p}{S}_r$ are geodesic curvatures of the curves belonging to the net considered and are defined by

$$\overset{p}{S}_r = \overset{p}{T}_k^p v^k.$$

In terms of the geodesic vector fields of the net, (2.13) reduces to

$$\overset{p}{C}_r^\alpha = \sum_{\sigma=1}^{n-m} w_{ik} v^i v^k n^\alpha + x_h^\alpha c_r^h + A_{ik}^h v^i v^k x_h^\alpha \quad (2.14)$$

where $\overset{i}{c}_r$ is the geodesic vector field, relative to GW_m , of the r -th family of the net defined by

$$\overset{i}{c}_r = \overset{p}{S}_r v^i.$$

The relations (2.6), (2.11), (2.14), allow us to state the following theorem:

Theorem 2.2. (i) A necessary and the sufficient condition the net $\delta = (v_1, v_2, \dots, v_n)$ in GW_m which is a chebyshev net of the first kind with respect to GW_n to be a chebyshev net of the first kind with respect to GW_m is that

$$\sum_{\sigma=1}^{n-m} w_{ik} v^i v^i = 0 \quad \text{and} \quad A_{ik}^h v^i v^k = 0$$

(ii) A necessary and the sufficient condition the net $\delta = (v_1, v_2, \dots, v_n)$ in GW_m which is a chebyshev net of the second kind with respect to GW_n to be a chebyshev net of the second kind with respect to GW_m is that

$$\sum_{\sigma=1}^{n-m} \Omega_k^i \underset{\textcircled{R}}{v_i} v^k = 0 \quad \text{and} \quad A_{hk}^i \underset{\textcircled{R}}{v_i} v^k = 0$$

(iii) A necessary and the sufficient condition the net $\delta = (v_1, v_2, \dots, v_n)$ in GW_m which is a geodesic net relative GW_n to be a geodesic net relative to GW_m is that

$$\sum_{\sigma=1}^{n-m} w_{ik} v_r^i v_r^k = 0 \quad \text{and} \quad A_{ik}^h v_r^i v_r^k = 0$$

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